

A High-Order Galerkin Solver for the Poisson Problem on the Surface of the Cubed Sphere

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Outline

- 1 Background
 - Spectral Element Method / Cubed Sphere
- 2 Poisson Solver
- 3 Results

Basic Premise

Goal: solve $-\nabla^2 u = f$ on the surface of a sphere.

Why?

- Eventually will solve global Shallow Water equations in vorticity-divergence form.
- Vorticity and divergence related to stream functions / velocity potentials via Laplace operator.

How?

- Spectral Element Method on the cubed sphere.

Spectral Element Method

- 1 Partition spatial domain into elements Ω^e .
- 2 Solve in **weak form**: multiply by a **test function** and integrate over each element.

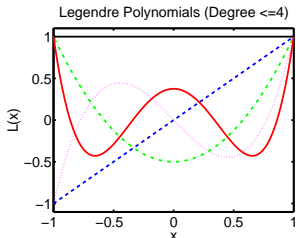
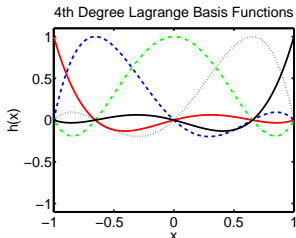
$$-\int_{\Omega^e} (\nabla^2 u) \phi d\Omega = \int_{\Omega^e} f \phi d\Omega$$

2D Basics

- Looking for an approximate solution $u_h(x, y) \approx u(x, y)$ in a finite-dimensional function space.
- Let $\mathcal{V}_h = \{v(x, y) : v(x, y) = p_n(x)q_n(y)\}$, where p_n and q_n are polynomials of degree $\leq n$.
 - Functions in \mathcal{V}_h must be continuous over element boundaries.
 - Both u_h and ϕ are in \mathcal{V}_h .
- Also need a quadrature rule for evaluating integrals: the **Gauss-Lobatto-Legendre** rule is used in both x and y
 \implies GLL nodes affinely mapped to elements define grid

More SEM Set-up (Spatial Discretization)

- Gaussian Quadrature: $\int_{-1}^1 \omega(x)p(x)dx = \sum_{i=0}^n w_i p(x_i)$
GLL: $\omega(x) \equiv 1$, $x_0 = -1$, $x_n = 1$
- Interpolation: two options for basis functions



Nodal expansion (Lagrange basis)

$$h_i(x_j) = \begin{cases} 1 & i = j \\ 0 & i \neq j \end{cases}$$

$$f(x) \approx \sum f(x_i)h_i(x)$$

Modal expansion (Legendre basis)

$$f(x) \approx \sum_i f_i L_i(x)$$

$$f_i = \int_{-1}^1 f(x)L_i(x)dx$$

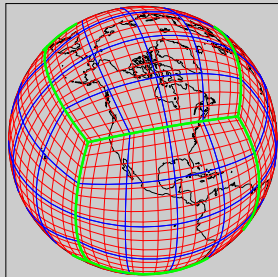
Cubed Sphere Basics

How can this methodology be extended to a spherical domain?

To use rectangular elements, turn to **cubed sphere**.

Cube Sphere Set-up

- 1 A cube is inscribed in a sphere
- 2 Points on the surface of the cube are projected onto the sphere
 - **Gnomic / central projection**
(ray from center to surface of sphere)
- 3 Each face is tiled with elements as in the 2D case



In the figure above, each face of the cube has a 4×4 element grid with a 6×6 GLL grid.

Mapping Vectors Between the Sphere and the Cube

For a sphere of radius a :

- On each face, the **metric tensor** g_{ij} is given by

$$g_{ij} = \frac{a^2}{\rho^4 \cos^2 x^1 \cos^2 x^2} \begin{bmatrix} 1 + \tan^2 x^1 & -\tan x^1 \tan x^2 \\ -\tan x^1 \tan x^2 & 1 + \tan^2 x^2 \end{bmatrix},$$

where $\rho = (1 + \tan^2 x^1 + \tan^2 x^2)^{1/2}$.

- Defining the matrix A by

$$A = a \begin{bmatrix} \cos \theta \partial \lambda / \partial x^1 & \cos \theta \partial \lambda / \partial x^2 \\ \partial \theta / \partial x^1 & \partial \theta / \partial x^2 \end{bmatrix},$$

where (x^1, x^2) are the cartesian coordinates on the face of the cube, it follows that $A^T A = g_{ij}$.

Laplacian Operator on the Cubed Sphere

In spherical coordinates, the Laplacian is given on the surface of a sphere by

$$\nabla^2 u = \nabla \cdot \nabla u = \frac{1}{a^2 \cos \theta} \frac{\partial}{\partial \theta} \left[\cos \theta \frac{\partial u}{\partial \theta} \right] + \frac{1}{a^2 \cos^2 \theta} \frac{\partial^2 u}{\partial \lambda^2}$$

On the surface of the cubed sphere, with $\nabla_g = (\partial/\partial x^1, \partial/\partial x^2)^T$,

$$\nabla^2 u = \frac{1}{\sqrt{g}} \nabla_g \cdot \left[\sqrt{g} A^{-1} A^{-T} \nabla_g u \right],$$

where $g = \det(g_{ij}) \implies \sqrt{g} = a^2 / (\rho^3 \cos^2 x^1 \cos^2 x^2)$.

Weak Form

(1) So the problem to solve is

$$-\frac{1}{\sqrt{g}} \nabla_g \cdot \left[\sqrt{g} A^{-1} A^{-T} \nabla_g u \right] = f.$$

(2) Or, slightly re-arranging terms,

$$-\nabla_g \cdot \left[\sqrt{g} A^{-1} A^{-T} \nabla_g u \right] = f \sqrt{g}.$$

(3) The first step is to cast in weak form:

$$-\int_{\Omega^e} \nabla_g \cdot \left[\sqrt{g} A^{-1} A^{-T} \nabla_g u \right] \phi d\Omega = \int_{\Omega^e} f \phi \sqrt{g} d\Omega.$$

(4) Integrating by parts simplifies the calculations:

$$\int_{\Omega^e} (A^{-T} \nabla_g u) \cdot (A^{-T} \nabla_g \phi) \sqrt{g} d\Omega = \int_{\Omega^e} f \phi \sqrt{g} d\Omega.$$

Quadrature

Letting $\phi(x^1, x^2) = h_p(x^1)h_q(x^2)$ for $p, q \in \{1, \dots, N\}$, and applying the GLL quadrature to the weak form, results in the linear system

$$K^e \mathbf{u}^e = M^e \mathbf{f}^e,$$

where \mathbf{u}^e and \mathbf{f}^e are vectors containing the nodal coefficients of u and f on Ω^e , respectively.

The solution must be continuous across element boundaries, and this is enforced by using [global assembly](#) to construct a global system: $K = \bigwedge_e K^e$, $M = \bigwedge_e M^e$ and the system

$$K\mathbf{u} = M\mathbf{f}$$

is solved using the [conjugate gradient](#) method.

Test Problem

If $u = \sin(\lambda) \cos(\theta) + C$ then $-\nabla^2 u = 2 \sin(\lambda) \cos(\theta) / a^2$. Working backwards, the test problem solved is

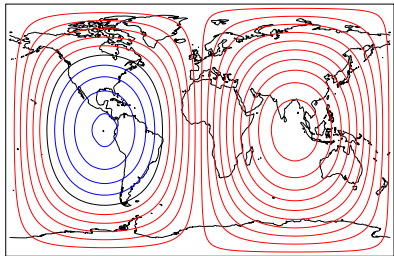
$$-\nabla^2 u = \frac{2 \sin(\lambda) \cos(\theta)}{a^2}$$

The numerical solution u_h is compared to the true solution $u = \sin(\lambda) \cos(\theta) + C$. The GLL quadrature rule is used to calculate the [relative L2 error](#):

$$\epsilon = \left(\frac{\int_{\Omega} (u - u_h)^2 \sqrt{g} d\Omega}{\int_{\Omega} u^2 \sqrt{g} d\Omega} \right)^{1/2}$$

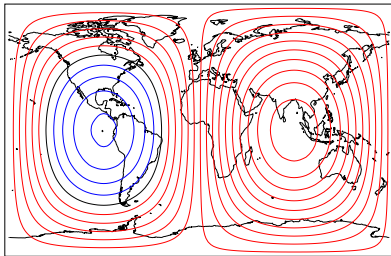
Contour Plots

Numerical Solution



Contour plot of u_h .

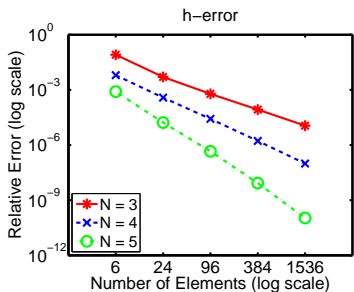
True Solution



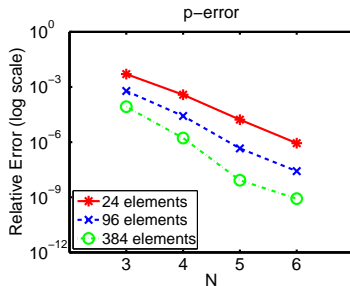
Contour plot of u .

For both plots, each face of the cube sphere had a 6×6 grid of elements and each element had a 4×4 GLL grid.

Error Plots



The *h-error* is measured by leaving the number of nodes per element constant but increasing the number of elements.



The *p-error* is measured by leaving the number of elements constant but increasing the number of nodes per element.

Future Work

- 1 **Parallelization**: this method is expensive, but fairly local so it should scale well.
- 2 **Preconditioning**: The conjugate gradient method is converging slowly for bigger grids / more elements; a diagonal preconditioner has been implemented but a better option may be needed.
- 3 **Shallow Water Model**: the work presented here, combined with an advection solver, will provide a high-order method for solving the shallow water equations (more at *PDEs on a Sphere '07*).